

**p** $I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$  Status: \*\*\***p MASS (atomic mass units u)**

The mass is known much more precisely in u (atomic mass units) than in MeV. See the next data block.

VALUE (u)	DOCUMENT ID	TECN	COMMENT
<b>1.007276466879±0.000000000091</b>	MOHR	16	RVUE 2014 CODATA value
• • • We do not use the following data for averages, fits, limits, etc. • • •			
1.007276466812±0.000000000090	MOHR	12	RVUE 2010 CODATA value
1.00727646677 ±0.00000000010	MOHR	08	RVUE 2006 CODATA value
1.00727646688 ±0.00000000013	MOHR	05	RVUE 2002 CODATA value
1.00727646688 ±0.00000000013	MOHR	99	RVUE 1998 CODATA value
1.007276470 ±0.000000012	COHEN	87	RVUE 1986 CODATA value

**p MASS (MeV)**

The mass is known much more precisely in u (atomic mass units) than in MeV. The conversion from u to MeV,  $1\text{ u} = 931.494\ 0054(57)\text{ MeV}/c^2$  (MOHR 16, the 2014 CODATA value), involves the relatively poorly known electronic charge.

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
<b>938.2720813±0.0000058</b>	MOHR	16	RVUE 2014 CODATA value
• • • We do not use the following data for averages, fits, limits, etc. • • •			
938.272046 ±0.000021	MOHR	12	RVUE 2010 CODATA value
938.272013 ±0.000023	MOHR	08	RVUE 2006 CODATA value
938.272029 ±0.0000080	MOHR	05	RVUE 2002 CODATA value
938.271998 ±0.0000038	MOHR	99	RVUE 1998 CODATA value
938.27231 ±0.00028	COHEN	87	RVUE 1986 CODATA value
938.2796 ±0.0027	COHEN	73	RVUE 1973 CODATA value

 **$|m_p - m_{\bar{p}}|/m_p$** 

A test of *CPT* invariance. Note that the comparison of the  $\bar{p}$  and  $p$  charge-to-mass ratio, given in the next data block, is much better determined.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<7 × 10 <sup>-10</sup>	90	<sup>1</sup> HORI	11	SPEC $\bar{p}e^-$ He atom
• • • We do not use the following data for averages, fits, limits, etc. • • •				
<2 × 10 <sup>-9</sup>	90	<sup>1</sup> HORI	06	SPEC $\bar{p}e^-$ He atom
<1.0 × 10 <sup>-8</sup>	90	<sup>1</sup> HORI	03	SPEC $\bar{p}e^-$ <sup>4</sup> He, $\bar{p}e^-$ <sup>3</sup> He
<6 × 10 <sup>-8</sup>	90	<sup>1</sup> HORI	01	SPEC $\bar{p}e^-$ He atom
<5 × 10 <sup>-7</sup>		<sup>2</sup> TORII	99	SPEC $\bar{p}e^-$ He atom

<sup>1</sup>HORI 01, HORI 03, HORI 06, and HORI 11 use the more-precisely-known constraint on the  $\bar{p}$  charge-to-mass ratio of GABRIELSE 99 (see below) to get their results. Their

results are not independent of the Hori 01, Hori 03, Hori 06, and Hori 11 values for  $|q_p + q_{\bar{p}}|/e$ , below.

<sup>2</sup> TORII 99 uses the more-precisely-known constraint on the  $\bar{p}$  charge-to-mass ratio of GABRIELSE 95 (see below) to get this result. This is not independent of the TORII 99 value for  $|q_p + q_{\bar{p}}|/e$ , below.

## $\bar{p}/p$ CHARGE-TO-MASS RATIO, $|q_{\bar{p}}/(m_{\bar{p}})|$

A test of *CPT* invariance. Listed here are measurements involving the *inertial* masses. For a discussion of what may be inferred about the ratio of  $\bar{p}$  and  $p$  *gravitational* masses, see ERICSON 90; they obtain an upper bound of  $10^{-6}$ – $10^{-7}$  for violation of the equivalence principle for  $\bar{p}$ 's.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.99999999991 ± 0.00000000009</b>	GABRIELSE 99	TRAP	Penning trap
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>			
1.0000000015 ± 0.0000000011	<sup>1</sup> GABRIELSE 95	TRAP	Penning trap
1.000000023 ± 0.000000042	<sup>2</sup> GABRIELSE 90	TRAP	Penning trap
<sup>1</sup> Equation (2) of GABRIELSE 95 should read $M(\bar{p})/M(p) = 0.999\ 999\ 9985$ (11) (G. Gabrielse, private communication).			
<sup>2</sup> GABRIELSE 90 also measures $m_{\bar{p}}/m_{e^-} = 1836.152660 \pm 0.000083$ and $m_p/m_{e^-} = 1836.152680 \pm 0.000088$ . Both are completely consistent with the 1986 CODATA (COHEN 87) value for $m_p/m_{e^-}$ of 1836.152701 ± 0.000037.			

$$(|\frac{q_{\bar{p}}}{m_{\bar{p}}} - \frac{q_p}{m_p}|)/\frac{q_p}{m_p}$$

A test of *CPT* invariance. Taken from the  $\bar{p}/p$  charge-to-mass ratio, above.

VALUE	DOCUMENT ID
<b>(−9±9) × 10<sup>−11</sup> OUR EVALUATION</b>	

$$|q_p + q_{\bar{p}}|/e$$

A test of *CPT* invariance. Note that the comparison of the  $\bar{p}$  and  $p$  charge-to-mass ratios given above is much better determined. See also a similar test involving the electron.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;7 × 10<sup>−10</sup></b>	90	<sup>1</sup> Hori 11	SPEC	$\bar{p}e^-$ He atom
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
<2 × 10 <sup>−9</sup>	90	<sup>1</sup> Hori 06	SPEC	$\bar{p}e^-$ He atom
<1.0 × 10 <sup>−8</sup>	90	<sup>1</sup> Hori 03	SPEC	$\bar{p}e^-$ <sup>4</sup> He, $\bar{p}e^-$ <sup>3</sup> He
<6 × 10 <sup>−8</sup>	90	<sup>1</sup> Hori 01	SPEC	$\bar{p}e^-$ He atom
<5 × 10 <sup>−7</sup>		<sup>2</sup> TORII 99	SPEC	$\bar{p}e^-$ He atom
<2 × 10 <sup>−5</sup>		<sup>3</sup> HUGHES 92	RVUE	

<sup>1</sup> Hori 01, Hori 03, Hori 06, and Hori 11 use the more-precisely-known constraint on the  $\bar{p}$  charge-to-mass ratio of GABRIELSE 99 (see above) to get their results. Their

results are not independent of the Hori 01, Hori 03, Hori 06, and Hori 11 values for  $|m_p - m_{\bar{p}}|/m_p$ , above.

<sup>2</sup>TORII 99 uses the more-precisely-known constraint on the  $\bar{p}$  charge-to-mass ratio of GABRIELSE 95 (see above) to get this result. This is not independent of the TORII 99 value for  $|m_p - m_{\bar{p}}|/m_p$ , above.

<sup>3</sup>HUGHES 92 uses recent measurements of Rydberg-energy and cyclotron-frequency ratios.

## $|q_p + q_e|/e$

See BRESSI 11 for a summary of experiments on the neutrality of matter.

See also “ $n$  CHARGE” in the neutron Listings.

VALUE	DOCUMENT ID	COMMENT
$<1 \times 10^{-21}$	<sup>1</sup> BRESSI 11	Neutrality of SF <sub>6</sub>
• • • We do not use the following data for averages, fits, limits, etc. • • •		
$<3.2 \times 10^{-20}$	<sup>2</sup> SENGUPTA 00	binary pulsar
$<0.8 \times 10^{-21}$	MARINELLI 84	Magnetic levitation
$<1.0 \times 10^{-21}$	<sup>1</sup> DYLLA 73	Neutrality of SF <sub>6</sub>

<sup>1</sup>BRESSI 11 uses the method of DYLLA 73 but finds serious errors in that experiment that greatly reduce its accuracy. The BRESSI 11 limit assumes that  $n \rightarrow p e^- \nu_e$  conserves charge. Thus the limit applies equally to the charge of the neutron.

<sup>2</sup>SENGUPTA 00 uses the difference between the observed rate of rotational energy loss by the binary pulsar PSR B1913+16 and the rate predicted by general relativity to set this limit. See the paper for assumptions.

## $p$ MAGNETIC MOMENT

See the “Note on Baryon Magnetic Moments” in the  $\Lambda$  Listings.

VALUE ( $\mu_N$ )	DOCUMENT ID	TECN	COMMENT
<b>2.7928473508 ± 0.0000000085</b>	MOHR 16	RVUE	2014 CODATA value
• • • We do not use the following data for averages, fits, limits, etc. • • •			
2.792847356 ± 0.000000023	MOHR 12	RVUE	2010 CODATA value
2.792847356 ± 0.000000023	MOHR 08	RVUE	2006 CODATA value
2.792847351 ± 0.000000028	MOHR 05	RVUE	2002 CODATA value
2.792847337 ± 0.000000029	MOHR 99	RVUE	1998 CODATA value
2.792847386 ± 0.000000063	COHEN 87	RVUE	1986 CODATA value
2.7928456 ± 0.0000011	COHEN 73	RVUE	1973 CODATA value

## $\bar{p}$ MAGNETIC MOMENT

A few early results have been omitted.

VALUE ( $\mu_N$ )	DOCUMENT ID	TECN	COMMENT
<b>-2.792845 ± 0.000012</b>	DISCIACCA 13	TRAP	Single $\bar{p}$ , Penning trap
• • • We do not use the following data for averages, fits, limits, etc. • • •			
-2.7862 ± 0.0083	PASK 09	CNTR	$\bar{p}$ He <sup>+</sup> hyperfine structure
-2.8005 ± 0.0090	KREISSL 88	CNTR	$\bar{p}$ <sup>208</sup> Pb 11 → 10 X-ray
-2.817 ± 0.048	ROBERTS 78	CNTR	

-2.791 ± 0.021

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75 CNTR Exotic atoms

 **$(\mu_p + \mu_{\bar{p}}) / \mu_p$** A test of *CPT* invariance.

VALUE (units $10^{-6}$ )	DOCUMENT ID	TECN	COMMENT
<b>0±5</b>	DISCIACCA	13	TRAP Single $\bar{p}$ , Penning trap

 **$p$  ELECTRIC DIPOLE MOMENT**A nonzero value is forbidden by both *T* invariance and *P* invariance.

VALUE ( $10^{-23}$ ecm)	EVTS	DOCUMENT ID	TECN	COMMENT
< <b>0.54</b>		1 DMITRIEV 03		Uses $^{199}\text{Hg}$ atom EDM
• • • We do not use the following data for averages, fits, limits, etc. • • •				
- 3.7 ± 6.3		CHO	89	NMR Ti F molecules
< 400		DZUBA	85	THEO Uses $^{129}\text{Xe}$ moment
130 ± 200		2 WILKENING	84	
900 ± 1400		3 WILKENING	84	
700 ± 900	1G	HARRISON	69	MBR Molecular beam

<sup>1</sup> DMITRIEV 03 calculates this limit from the limit on the electric dipole moment of the  $^{199}\text{Hg}$  atom.

<sup>2</sup> This WILKENING 84 value includes a finite-size effect and a magnetic effect.

<sup>3</sup> This WILKENING 84 value is more cautious than the other and excludes the finite-size effect, which relies on uncertain nuclear integrals.

 **$p$  ELECTRIC POLARIZABILITY  $\alpha_p$** 

For a very complete review of the “polarizability of the nucleon and Compton scattering,” see SCHUMACHER 05. His recommended values for the proton are  $\alpha_p = (12.0 \pm 0.6) \times 10^{-4} \text{ fm}^3$  and  $\beta_p = (1.9 \mp 0.6) \times 10^{-4} \text{ fm}^3$ , almost exactly our averages.

VALUE ( $10^{-4} \text{ fm}^3$ )	DOCUMENT ID	TECN	COMMENT
<b>11.2 ± 0.4 OUR AVERAGE</b>			
10.65 ± 0.35 ± 0.36	MCGOVERN 13	RVUE	$\chi$ EFT + Compton scattering
12.1 ± 1.1 ± 0.5	<sup>1</sup> BEANE 03		EFT + $\gamma p$
11.82 ± 0.98 ± 0.52	2 BLANPIED 01	LEGS	$p(\vec{\gamma}, \gamma)$ , $p(\vec{\gamma}, \pi^0)$ , $p(\vec{\gamma}, \pi^+)$
11.9 ± 0.5 ± 1.3	<sup>3</sup> OLMSDEL... 01	CNTR	$\gamma p$ Compton scattering
12.1 ± 0.8 ± 0.5	<sup>4</sup> MACGIBBON 95	RVUE	global average
• • • We do not use the following data for averages, fits, limits, etc. • • •			
11.7 ± 0.8 ± 0.7	<sup>5</sup> BARANOV 01	RVUE	Global average
12.5 ± 0.6 ± 0.9	MACGIBBON 95	CNTR	$\gamma p$ Compton scattering
9.8 ± 0.4 ± 1.1	HALLIN 93	CNTR	$\gamma p$ Compton scattering
10.62 ± 1.25 ± 1.07	ZIEGER 92	CNTR	$\gamma p$ Compton scattering
10.9 ± 2.2 ± 1.3	<sup>6</sup> FEDERSPIEL 91	CNTR	$\gamma p$ Compton scattering

- <sup>1</sup> BEANE 03 uses effective field theory and low-energy  $\gamma p$  and  $\gamma d$  Compton-scattering data. It also gets for the isoscalar polarizabilities (see the erratum)  $\alpha_N = (13.0 \pm 1.9^{+3.9}_{-1.5}) \times 10^{-4} \text{ fm}^3$  and  $\beta_N = (-1.8 \pm 1.9^{+2.1}_{-0.9}) \times 10^{-4} \text{ fm}^3$ .
- <sup>2</sup> BLANPIED 01 gives  $\alpha_p + \beta_p$  and  $\alpha_p - \beta_p$ . The separate  $\alpha_p$  and  $\beta_p$  are provided to us by A. Sandorfi. The first error above is statistics plus systematics; the second is from the model.
- <sup>3</sup> This OLMSDELEON 01 result uses the TAPS data alone, and does not use the (re-evaluated) sum-rule constraint that  $\alpha + \beta = (13.8 \pm 0.4) \times 10^{-4} \text{ fm}^3$ . See the paper for a discussion.
- <sup>4</sup> MACGIBBON 95 combine the results of ZIEGER 92, FEDERSPIEL 91, and their own experiment to get a “global average” in which model errors and systematic errors are treated in a consistent way. See MACGIBBON 95 for a discussion.
- <sup>5</sup> BARANOV 01 combines the results of 10 experiments from 1958 through 1995 to get a global average that takes into account both systematic and model errors and does not use the theoretical constraint on the sum  $\alpha_p + \beta_p$ .
- <sup>6</sup> FEDERSPIEL 91 obtains for the (static) electric polarizability  $\alpha_p$ , defined in terms of the induced electric dipole moment by  $\mathbf{D} = 4\pi\epsilon_0\alpha_p\mathbf{E}$ , the value  $(7.0 \pm 2.2 \pm 1.3) \times 10^{-4} \text{ fm}^3$ .

## $p$ MAGNETIC POLARIZABILITY $\beta_p$

The electric and magnetic polarizabilities are subject to a dispersion sum-rule constraint  $\overline{\alpha} + \overline{\beta} = (14.2 \pm 0.5) \times 10^{-4} \text{ fm}^3$ . Errors here are anticorrelated with those on  $\overline{\alpha}_p$  due to this constraint.

VALUE ( $10^{-4} \text{ fm}^3$ )	DOCUMENT ID	TECN	COMMENT
<b>2.5 <math>\pm 0.4</math> OUR AVERAGE</b>	Error includes scale factor of 1.2.		
3.15 $\pm 0.35 \pm 0.36$	MCGOVERN 13	RVUE	$\chi$ EFT + Compton scattering
3.4 $\pm 1.1 \pm 0.1$	<sup>1</sup> BEANE 03		EFT + $\gamma p$
1.43 $\pm 0.98^{+0.52}_{-0.98}$	<sup>2</sup> BLANPIED 01	LEGS	$p(\vec{\gamma}, \gamma)$ , $p(\vec{\gamma}, \pi^0)$ , $p(\vec{\gamma}, \pi^+)$
1.2 $\pm 0.7 \pm 0.5$	<sup>3</sup> OLMSDEL... 01	CNTR	$\gamma p$ Compton scattering
2.1 $\pm 0.8 \pm 0.5$	<sup>4</sup> MACGIBBON 95	RVUE	global average
• • • We do not use the following data for averages, fits, limits, etc. • • •			
2.3 $\pm 0.9 \pm 0.7$	<sup>5</sup> BARANOV 01	RVUE	Global average
1.7 $\pm 0.6 \pm 0.9$	MACGIBBON 95	CNTR	$\gamma p$ Compton scattering
4.4 $\pm 0.4 \pm 1.1$	HALLIN 93	CNTR	$\gamma p$ Compton scattering
3.58 $\pm 1.19 \pm 1.03$	ZIEGER 92	CNTR	$\gamma p$ Compton scattering
3.3 $\pm 2.2 \pm 1.3$	FEDERSPIEL 91	CNTR	$\gamma p$ Compton scattering

- <sup>1</sup> BEANE 03 uses effective field theory and low-energy  $\gamma p$  and  $\gamma d$  Compton-scattering data. It also gets for the isoscalar polarizabilities (see the erratum)  $\alpha_N = (13.0 \pm 1.9^{+3.9}_{-1.5}) \times 10^{-4} \text{ fm}^3$  and  $\beta_N = (-1.8 \pm 1.9^{+2.1}_{-0.9}) \times 10^{-4} \text{ fm}^3$ .
- <sup>2</sup> BLANPIED 01 gives  $\alpha_p + \beta_p$  and  $\alpha_p - \beta_p$ . The separate  $\alpha_p$  and  $\beta_p$  are provided to us by A. Sandorfi. The first error above is statistics plus systematics; the second is from the model.
- <sup>3</sup> This OLMSDELEON 01 result uses the TAPS data alone, and does not use the (re-evaluated) sum-rule constraint that  $\alpha + \beta = (13.8 \pm 0.4) \times 10^{-4} \text{ fm}^3$ . See the paper for a discussion.
- <sup>4</sup> MACGIBBON 95 combine the results of ZIEGER 92, FEDERSPIEL 91, and their own experiment to get a “global average” in which model errors and systematic errors are treated in a consistent way. See MACGIBBON 95 for a discussion.

<sup>5</sup> BARANOV 01 combines the results of 10 experiments from 1958 through 1995 to get a global average that takes into account both systematic and model errors and does not use the theoretical constraint on the sum  $\alpha_p + \beta_p$ .

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## $p$ CHARGE RADIUS

This is the rms electric charge radius,  $\sqrt{\langle r_E^2 \rangle}$ .

Most measurements of the radius of the proton involve electron-proton interactions, and most of the more recent values agree with one another. The most precise of these is  $r_p = 0.879(8)$  fm (BERNAUER 10). The CODATA 14 value (MOHR 16), obtained from the electronic results, is 0.8751(61). Compared to this CODATA value, however, a measurement using muonic hydrogen finds  $r_p = 0.84087(39)$  fm (ANTOGNINI 13), which is 16 times more precise but differs by 5.6 standard deviations (using the CODATA 14 error).

Since POHL 10 (the first  $\mu p$  result), there has been a lot of discussion about the disagreement, especially concerning the modeling of muonic hydrogen. Here is an incomplete list of papers: DERUJULA 10, CLOET 11, DISTLER 11, DERUJULA 11, ARRINGTON 11, BERNAUER 11, HILL 11, LORENZ 14, KARSHENBOIM 14A, and PESET 15.

Until the difference between the  $e p$  and  $\mu p$  values is understood, it does not make sense to average the values together. For the present, we give both values. It is up to workers in this field to solve this puzzle.

See our 2014 edition (Chinese Physics C **38** 070001 (2014)) for values published before 2003.

VALUE (fm)	DOCUMENT ID	TECN	COMMENT
<b>0.8751 ± 0.0061</b>	MOHR	16	RVUE 2014 CODATA value
<b>0.84087±0.00026±0.00029</b>	ANTOGNINI	13	LASR $\mu p$ -atom Lamb shift
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.895 ± 0.014 ± 0.014	<sup>1</sup> LEE	15	SPEC Just 2010 Mainz data
0.916 ± 0.024	LEE	15	SPEC World data, no Mainz
0.8775 ± 0.0051	MOHR	12	RVUE 2010 CODATA, $e p$ data
0.875 ± 0.008 ± 0.006	ZHAN	11	SPEC Recoil polarimetry
0.879 ± 0.005 ± 0.006	BERNAUER	10	SPEC $e p \rightarrow e p$ form factor
0.912 ± 0.009 ± 0.007	BORISYUK	10	reanalyzes old $e p$ data
0.871 ± 0.009 ± 0.003	HILL	10	$z$ -expansion reanalysis
0.84184±0.00036±0.00056	POHL	10	LASR See ANTOGNINI 13
0.8768 ± 0.0069	MOHR	08	RVUE 2006 CODATA value
0.844 $^{+0.008}_{-0.004}$	BELUSHKIN	07	Dispersion analysis
0.897 ± 0.018	BLUNDEN	05	SICK 03 + $2\gamma$ correction
0.8750 ± 0.0068	MOHR	05	RVUE 2002 CODATA value
0.895 ± 0.010 ± 0.013	SICK	03	$e p \rightarrow e p$ reanalysis

<sup>1</sup> Authors also provide values for combinations of all available data.

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## **p MAGNETIC RADIUS**

This is the rms magnetic radius,  $\sqrt{\langle r_M^2 \rangle}$ .

VALUE (fm)	DOCUMENT ID	TECN	COMMENT
<b>0.776±0.034±0.017</b>	<sup>1</sup> LEE	15	SPEC Just 2010 Mainz data
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>			
0.914±0.035	LEE	15	SPEC World data, no Mainz
0.87 ± 0.02	EPSTEIN	14	Using $e p$ , $e n$ , $\pi\pi$ data
0.867±0.009±0.018	ZHAN	11	SPEC Recoil polarimetry
0.777±0.013±0.010	BERNAUER	10	SPEC $e p \rightarrow e p$ form factor
0.876±0.010±0.016	BORISYUK	10	Reanalyzes old $e p \rightarrow e p$ data
0.854±0.005	BELUSHKIN	07	Dispersion analysis

<sup>1</sup> Authors also provide values for a combination of all available data.

## **p MEAN LIFE**

A test of baryon conservation. See the “*p* Partial Mean Lives” section below for limits for identified final states. The limits here are to “anything” or are for “disappearance” modes of a bound proton (*p*) or (*n*). See also the  $3\nu$  modes in the “Partial Mean Lives” section. Table 1 of BACK 03 is a nice summary.

LIMIT (years)	PARTICLE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&gt;5.8 × 10<sup>29</sup></b>	<b>n</b>	90	<sup>1</sup> ARAKI	06	KLND $n \rightarrow$ invisible
<b>&gt;2.1 × 10<sup>29</sup></b>	<b>p</b>	90	<sup>2</sup> AHMED	04	SNO $p \rightarrow$ invisible
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>					
>1.9 × 10 <sup>29</sup>	<i>n</i>	90	<sup>2</sup> AHMED	04	SNO $n \rightarrow$ invisible
>1.8 × 10 <sup>25</sup>	<i>n</i>	90	<sup>3</sup> BACK	03	BORX
>1.1 × 10 <sup>26</sup>	<i>p</i>	90	<sup>3</sup> BACK	03	BORX
>3.5 × 10 <sup>28</sup>	<i>p</i>	90	<sup>4</sup> ZDESENKO	03	$p \rightarrow$ invisible
>1 × 10 <sup>28</sup>	<i>p</i>	90	<sup>5</sup> AHMAD	02	SNO $p \rightarrow$ invisible
>4 × 10 <sup>23</sup>	<i>p</i>	95	TRETYAK	01	$d \rightarrow n + ?$
>1.9 × 10 <sup>24</sup>	<i>p</i>	90	<sup>6</sup> BERNABEI	00B	DAMA
>1.6 × 10 <sup>25</sup>	<i>p, n</i>		<sup>7,8</sup> EVANS	77	
>3 × 10 <sup>23</sup>	<i>p</i>		<sup>8</sup> DIX	70	CNTR
>3 × 10 <sup>23</sup>	<i>p, n</i>		<sup>8,9</sup> FLEROV	58	

<sup>1</sup> ARAKI 06 looks for signs of de-excitation of the residual nucleus after disappearance of a neutron from the *s* shell of  $^{12}\text{C}$ .

<sup>2</sup> AHMED 04 looks for  $\gamma$  rays from the de-excitation of a residual  $^{15}\text{O}^*$  or  $^{15}\text{N}^*$  following the disappearance of a neutron or proton in  $^{16}\text{O}$ .

<sup>3</sup> BACK 03 looks for decays of unstable nuclides left after *N* decays of parent  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$  nuclei. These are “invisible channel” limits.

<sup>4</sup> ZDESENKO 03 gets this limit on proton disappearance in deuterium by analyzing SNO data in AHMAD 02.

<sup>5</sup> AHMAD 02 (see its footnote 7) looks for neutrons left behind after the disappearance of the proton in deuterons.

<sup>6</sup> BERNABEI 00B looks for the decay of a  $^{128}_{53}\text{I}$  nucleus following the disappearance of a proton in the otherwise-stable  $^{129}_{54}\text{Xe}$  nucleus.

<sup>7</sup> EVANS 77 looks for the daughter nuclide  $^{129}\text{Xe}$  from possible  $^{130}\text{Te}$  decays in ancient Te ore samples.

<sup>8</sup> This mean-life limit has been obtained from a half-life limit by dividing the latter by  $\ln(2) = 0.693$ .

<sup>9</sup> FLEROV 58 looks for the spontaneous fission of a  $^{232}\text{Th}$  nucleus after the disappearance of one of its nucleons.

## $\bar{p}$ MEAN LIFE

Of the two astrophysical limits here, that of GEER 00D involves considerably more refinements in its modeling. The other limits come from direct observations of stored antiprotons. See also “ $\bar{p}$  Partial Mean Lives” after “ $p$  Partial Mean Lives,” below, for exclusive-mode limits. The best (lifetime/branching fraction) limit there is  $7 \times 10^5$  years, for  $\bar{p} \rightarrow e^- \gamma$ . We advance only the exclusive-mode limits to our Summary Tables.

LIMIT (years)	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>					
$>8 \times 10^5$	90		<sup>1</sup> GEER	00D	$\bar{p}/p$ ratio, cosmic rays
$>0.28$			GABRIELSE	90	TRAP Penning trap
$>0.08$	90	1	BELL	79	CNTR Storage ring
$>1 \times 10^7$			GOLDEN	79	SPEC $\bar{p}/p$ ratio, cosmic rays
$>3.7 \times 10^{-3}$			BREGMAN	78	CNTR Storage ring

<sup>1</sup> GEER 00D uses agreement between a model of galactic  $\bar{p}$  production and propagation and the observed  $\bar{p}/p$  cosmic-ray spectrum to set this limit.

## $p$ DECAY MODES

See the “Note on Nucleon Decay” in our 1994 edition (Phys. Rev. **D50**, 1173) for a short review.

The “partial mean life” limits tabulated here are the limits on  $\tau/B_i$ , where  $\tau$  is the total mean life and  $B_i$  is the branching fraction for the mode in question. For  $N$  decays,  $p$  and  $n$  indicate proton and neutron partial lifetimes.

Mode	Partial mean life ( $10^{30}$ years)	Confidence level
<b>Antilepton + meson</b>		
$\tau_1 \quad N \rightarrow e^+ \pi$	$> 2000 \text{ (} n \text{), } > 8200 \text{ (} p \text{)}$	90%
$\tau_2 \quad N \rightarrow \mu^+ \pi$	$> 1000 \text{ (} n \text{), } > 6600 \text{ (} p \text{)}$	90%
$\tau_3 \quad N \rightarrow \nu \pi$	$> 1100 \text{ (} n \text{), } > 390 \text{ (} p \text{)}$	90%
$\tau_4 \quad p \rightarrow e^+ \eta$	$> 4200$	90%
$\tau_5 \quad p \rightarrow \mu^+ \eta$	$> 1300$	90%
$\tau_6 \quad n \rightarrow \nu \eta$	$> 158$	90%
$\tau_7 \quad N \rightarrow e^+ \rho$	$> 217 \text{ (} n \text{), } > 710 \text{ (} p \text{)}$	90%
$\tau_8 \quad N \rightarrow \mu^+ \rho$	$> 228 \text{ (} n \text{), } > 160 \text{ (} p \text{)}$	90%
$\tau_9 \quad N \rightarrow \nu \rho$	$> 19 \text{ (} n \text{), } > 162 \text{ (} p \text{)}$	90%
$\tau_{10} \quad p \rightarrow e^+ \omega$	$> 320$	90%
$\tau_{11} \quad p \rightarrow \mu^+ \omega$	$> 780$	90%

$\tau_{12}$	$n \rightarrow \nu \omega$	$> 108$	90%
$\tau_{13}$	$N \rightarrow e^+ K$	$> 17 (n), > 1000 (p)$	90%
$\tau_{14}$	$p \rightarrow e^+ K_S^0$		
$\tau_{15}$	$p \rightarrow e^+ K_L^0$		
$\tau_{16}$	$N \rightarrow \mu^+ K$	$> 26 (n), > 1600 (p)$	90%
$\tau_{17}$	$p \rightarrow \mu^+ K_S^0$		
$\tau_{18}$	$p \rightarrow \mu^+ K_L^0$		
$\tau_{19}$	$N \rightarrow \nu K$	$> 86 (n), > 5900 (p)$	90%
$\tau_{20}$	$n \rightarrow \nu K_S^0$	$> 260$	90%
$\tau_{21}$	$p \rightarrow e^+ K^*(892)^0$	$> 84$	90%
$\tau_{22}$	$N \rightarrow \nu K^*(892)$	$> 78 (n), > 51 (p)$	90%

**Antilepton + mesons**

$\tau_{23}$	$p \rightarrow e^+ \pi^+ \pi^-$	$> 82$	90%
$\tau_{24}$	$p \rightarrow e^+ \pi^0 \pi^0$	$> 147$	90%
$\tau_{25}$	$n \rightarrow e^+ \pi^- \pi^0$	$> 52$	90%
$\tau_{26}$	$p \rightarrow \mu^+ \pi^+ \pi^-$	$> 133$	90%
$\tau_{27}$	$p \rightarrow \mu^+ \pi^0 \pi^0$	$> 101$	90%
$\tau_{28}$	$n \rightarrow \mu^+ \pi^- \pi^0$	$> 74$	90%
$\tau_{29}$	$n \rightarrow e^+ K^0 \pi^-$	$> 18$	90%

**Lepton + meson**

$\tau_{30}$	$n \rightarrow e^- \pi^+$	$> 65$	90%
$\tau_{31}$	$n \rightarrow \mu^- \pi^+$	$> 49$	90%
$\tau_{32}$	$n \rightarrow e^- \rho^+$	$> 62$	90%
$\tau_{33}$	$n \rightarrow \mu^- \rho^+$	$> 7$	90%
$\tau_{34}$	$n \rightarrow e^- K^+$	$> 32$	90%
$\tau_{35}$	$n \rightarrow \mu^- K^+$	$> 57$	90%

**Lepton + mesons**

$\tau_{36}$	$p \rightarrow e^- \pi^+ \pi^+$	$> 30$	90%
$\tau_{37}$	$n \rightarrow e^- \pi^+ \pi^0$	$> 29$	90%
$\tau_{38}$	$p \rightarrow \mu^- \pi^+ \pi^+$	$> 17$	90%
$\tau_{39}$	$n \rightarrow \mu^- \pi^+ \pi^0$	$> 34$	90%
$\tau_{40}$	$p \rightarrow e^- \pi^+ K^+$	$> 75$	90%
$\tau_{41}$	$p \rightarrow \mu^- \pi^+ K^+$	$> 245$	90%

**Antilepton + photon(s)**

$\tau_{42}$	$p \rightarrow e^+ \gamma$	$> 670$	90%
$\tau_{43}$	$p \rightarrow \mu^+ \gamma$	$> 478$	90%
$\tau_{44}$	$n \rightarrow \nu \gamma$	$> 550$	90%
$\tau_{45}$	$p \rightarrow e^+ \gamma \gamma$	$> 100$	90%
$\tau_{46}$	$n \rightarrow \nu \gamma \gamma$	$> 219$	90%

**Antilepton + single massless**

$\tau_{47}$	$p \rightarrow e^+ X$	$> 790$	90%
$\tau_{48}$	$p \rightarrow \mu^+ X$	$> 410$	90%

**Three (or more) leptons**

$\tau_{49}$	$p \rightarrow e^+ e^+ e^-$	$> 793$	90%
$\tau_{50}$	$p \rightarrow e^+ \mu^+ \mu^-$	$> 359$	90%
$\tau_{51}$	$p \rightarrow e^+ \nu \nu$	$> 170$	90%
$\tau_{52}$	$n \rightarrow e^+ e^- \nu$	$> 257$	90%
$\tau_{53}$	$n \rightarrow \mu^+ e^- \nu$	$> 83$	90%
$\tau_{54}$	$n \rightarrow \mu^+ \mu^- \nu$	$> 79$	90%
$\tau_{55}$	$p \rightarrow \mu^+ e^+ e^-$	$> 529$	90%
$\tau_{56}$	$p \rightarrow \mu^+ \mu^+ \mu^-$	$> 675$	90%
$\tau_{57}$	$p \rightarrow \mu^+ \nu \nu$	$> 220$	90%
$\tau_{58}$	$p \rightarrow e^- \mu^+ \mu^+$	$> 6$	90%
$\tau_{59}$	$n \rightarrow 3\nu$	$> 5 \times 10^{-4}$	90%
$\tau_{60}$	$n \rightarrow 5\nu$		

**Inclusive modes**

$\tau_{61}$	$N \rightarrow e^+ \text{anything}$	$> 0.6 (n, p)$	90%
$\tau_{62}$	$N \rightarrow \mu^+ \text{anything}$	$> 12 (n, p)$	90%
$\tau_{63}$	$N \rightarrow \nu \text{anything}$		
$\tau_{64}$	$N \rightarrow e^+ \pi^0 \text{anything}$	$> 0.6 (n, p)$	90%
$\tau_{65}$	$N \rightarrow 2 \text{ bodies, } \nu\text{-free}$		

 **$\Delta B = 2$  dinucleon modes**

The following are lifetime limits per iron nucleus.

$\tau_{66}$	$pp \rightarrow \pi^+ \pi^+$	$> 72.2$	90%
$\tau_{67}$	$pn \rightarrow \pi^+ \pi^0$	$> 170$	90%
$\tau_{68}$	$nn \rightarrow \pi^+ \pi^-$	$> 0.7$	90%
$\tau_{69}$	$nn \rightarrow \pi^0 \pi^0$	$> 404$	90%
$\tau_{70}$	$pp \rightarrow K^+ K^+$	$> 170$	90%
$\tau_{71}$	$pp \rightarrow e^+ e^+$	$> 5.8$	90%
$\tau_{72}$	$pp \rightarrow e^+ \mu^+$	$> 3.6$	90%
$\tau_{73}$	$pp \rightarrow \mu^+ \mu^+$	$> 1.7$	90%
$\tau_{74}$	$pn \rightarrow e^+ \bar{\nu}$	$> 260$	90%
$\tau_{75}$	$pn \rightarrow \mu^+ \bar{\nu}$	$> 200$	90%
$\tau_{76}$	$pn \rightarrow \tau^+ \bar{\nu}_\tau$	$> 29$	90%
$\tau_{77}$	$nn \rightarrow \nu_e \bar{\nu}_e$	$> 1.4$	90%
$\tau_{78}$	$nn \rightarrow \nu_\mu \bar{\nu}_\mu$	$> 1.4$	90%
$\tau_{79}$	$pn \rightarrow \text{invisible}$	$> 2.1 \times 10^{-5}$	90%
$\tau_{80}$	$pp \rightarrow \text{invisible}$	$> 5 \times 10^{-5}$	90%

**$\bar{p}$  DECAY MODES**

Mode		Partial mean life (years)	Confidence level
$\tau_{81}$	$\bar{p} \rightarrow e^- \gamma$	$> 7 \times 10^5$	90%
$\tau_{82}$	$\bar{p} \rightarrow \mu^- \gamma$	$> 5 \times 10^4$	90%
$\tau_{83}$	$\bar{p} \rightarrow e^- \pi^0$	$> 4 \times 10^5$	90%
$\tau_{84}$	$\bar{p} \rightarrow \mu^- \pi^0$	$> 5 \times 10^4$	90%
$\tau_{85}$	$\bar{p} \rightarrow e^- \eta$	$> 2 \times 10^4$	90%
$\tau_{86}$	$\bar{p} \rightarrow \mu^- \eta$	$> 8 \times 10^3$	90%
$\tau_{87}$	$\bar{p} \rightarrow e^- K_S^0$	$> 900$	90%
$\tau_{88}$	$\bar{p} \rightarrow \mu^- K_S^0$	$> 4 \times 10^3$	90%
$\tau_{89}$	$\bar{p} \rightarrow e^- K_L^0$	$> 9 \times 10^3$	90%
$\tau_{90}$	$\bar{p} \rightarrow \mu^- K_L^0$	$> 7 \times 10^3$	90%
$\tau_{91}$	$\bar{p} \rightarrow e^- \gamma\gamma$	$> 2 \times 10^4$	90%
$\tau_{92}$	$\bar{p} \rightarrow \mu^- \gamma\gamma$	$> 2 \times 10^4$	90%
$\tau_{93}$	$\bar{p} \rightarrow e^- \omega$	$> 200$	90%

 **$p$  PARTIAL MEAN LIVES**

The “partial mean life” limits tabulated here are the limits on  $\tau/B_i$ , where  $\tau$  is the total mean life for the proton and  $B_i$  is the branching fraction for the mode in question.

Decaying particle:  $p$  = proton,  $n$  = bound neutron. The same event may appear under more than one partial decay mode. Background estimates may be accurate to a factor of two.

**— Antilepton + meson —** **$\tau(N \rightarrow e^+ \pi)$**  **$\tau_1$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;2000</b>	<b><math>n</math></b>	<b>90</b>	<b>0</b>	<b>0.27</b>	NISHINO	12
<b>&gt;8200</b>	<b><math>p</math></b>	<b>90</b>	<b>0</b>	<b>0.3</b>	NISHINO	09

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 540	$p$	90	0	0.2	MCGREW	99	IMB3
> 158	$n$	90	3	5	MCGREW	99	IMB3
>1600	$p$	90	0	0.1	SHIOZAWA	98	SKAM
> 70	$p$	90	0	0.5	BERGER	91	FREJ
> 70	$n$	90	0	$\leq 0.1$	BERGER	91	FREJ
> 550	$p$	90	0	0.7	<sup>1</sup> BECKER-SZ...	90	IMB3
> 260	$p$	90	0	$<0.04$	HIRATA	89C	KAMI
> 130	$n$	90	0	$<0.2$	HIRATA	89C	KAMI
> 310	$p$	90	0	0.6	SEIDEL	88	IMB
> 100	$n$	90	0	1.6	SEIDEL	88	IMB
> 1.3	$n$	90	0		BARTEL	87	SOUD
> 1.3	$p$	90	0		BARTEL	87	SOUD

> 250	<i>p</i>	90	0 0.3	HAINES	86	IMB
> 31	<i>n</i>	90	8 9	HAINES	86	IMB
> 64	<i>p</i>	90	0 <0.4	ARISAKA	85	KAMI
> 26	<i>n</i>	90	0 <0.7	ARISAKA	85	KAMI
> 82	<i>p</i> (free)	90	0 0.2	BLEWITT	85	IMB
> 250	<i>p</i>	90	0 0.2	BLEWITT	85	IMB
> 25	<i>n</i>	90	4 4	PARK	85	IMB
> 15	<i>p, n</i>	90	0	BATTISTONI	84	NUSX
> 0.5	<i>p</i>	90	1 0.3	<sup>2</sup> BARTEL	83	SOUD
> 0.5	<i>n</i>	90	1 0.3	<sup>2</sup> BARTEL	83	SOUD
> 5.8	<i>p</i>	90	2	<sup>3</sup> KRISHNA...	82	KOLR
> 5.8	<i>n</i>	90	2	<sup>3</sup> KRISHNA...	82	KOLR
> 0.1	<i>n</i>	90		<sup>4</sup> GURR	67	CNTR

<sup>1</sup> This BECKER-SZENDY 90 result includes data from SEIDEL 88.<sup>2</sup> Limit based on zero events.<sup>3</sup> We have calculated 90% CL limit from 1 confined event.<sup>4</sup> We have converted half-life to 90% CL mean life. $\tau(N \rightarrow \mu^+ \pi^-)$  $\tau_2$ 

LIMIT ( $10^{-30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>1000	<i>n</i>	<b>90</b>	<b>1</b>	<b>0.43</b>	NISHINO	12
>6600	<i>p</i>	<b>90</b>	<b>0</b>	<b>0.3</b>	NISHINO	09

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 473	<i>p</i>	90	0 0.6	MCGREW	99	IMB3
> 90	<i>n</i>	90	1 1.9	MCGREW	99	IMB3
> 81	<i>p</i>	90	0 0.2	BERGER	91	FREJ
> 35	<i>n</i>	90	1 1.0	BERGER	91	FREJ
> 230	<i>p</i>	90	0 <0.07	HIRATA	89C	KAMI
> 100	<i>n</i>	90	0 <0.2	HIRATA	89C	KAMI
> 270	<i>p</i>	90	0 0.5	SEIDEL	88	IMB
> 63	<i>n</i>	90	0 0.5	SEIDEL	88	IMB
> 76	<i>p</i>	90	2 1	HAINES	86	IMB
> 23	<i>n</i>	90	8 7	HAINES	86	IMB
> 46	<i>p</i>	90	0 <0.7	ARISAKA	85	KAMI
> 20	<i>n</i>	90	0 <0.4	ARISAKA	85	KAMI
> 59	<i>p</i> (free)	90	0 0.2	BLEWITT	85	IMB
> 100	<i>p</i>	90	1 0.4	BLEWITT	85	IMB
> 38	<i>n</i>	90	1 4	PARK	85	IMB
> 10	<i>p, n</i>	90	0	BATTISTONI	84	NUSX
> 1.3	<i>p, n</i>	90	0	ALEKSEEV	81	BAKS

 $\tau(N \rightarrow \nu\pi)$  $\tau_3$ 

LIMIT ( $10^{-30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
> 390	<i>p</i>	<b>90</b>	<b>52.8</b>		ABE	14E
>1100	<i>n</i>	<b>90</b>	<b>19.1</b>		ABE	14E

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 16	<i>p</i>	90	6	6.7	WALL	00B	SOU2
> 39	<i>n</i>	90	4	3.8	WALL	00B	SOU2
> 10	<i>p</i>	90	15	20.3	MCGREW	99	IMB3
> 112	<i>n</i>	90	6	6.6	MCGREW	99	IMB3
> 13	<i>n</i>	90	1	1.2	BERGER	89	FREJ
> 10	<i>p</i>	90	11	14	BERGER	89	FREJ
> 25	<i>p</i>	90	32	32.8	<sup>1</sup> HIRATA	89C	KAMI
> 100	<i>n</i>	90	1	3	HIRATA	89C	KAMI
> 6	<i>n</i>	90	73	60	HAINES	86	IMB
> 2	<i>p</i>	90	16	13	KAJITA	86	KAMI
> 40	<i>n</i>	90	0	1	KAJITA	86	KAMI
> 7	<i>n</i>	90	28	19	PARK	85	IMB
> 7	<i>n</i>	90	0		BATTISTONI	84	NUSX
> 2	<i>p</i>	90	≤ 3		BATTISTONI	84	NUSX
> 5.8	<i>p</i>	90	1		<sup>2</sup> KRISHNA...	82	KOLR
> 0.3	<i>p</i>	90	2		<sup>3</sup> CHERRY	81	HOME
> 0.1	<i>p</i>	90			<sup>4</sup> GURR	67	CNTR

<sup>1</sup> In estimating the background, this HIRATA 89C limit (as opposed to the later limits of WALL 00B and MCGREW 99) does not take into account present understanding that the flux of  $\nu_\mu$  originating in the upper atmosphere is depleted. Doing so would reduce the background and thus also would reduce the limit here.

<sup>2</sup> We have calculated 90% CL limit from 1 confined event.

<sup>3</sup> We have converted 2 possible events to 90% CL limit.

<sup>4</sup> We have converted half-life to 90% CL mean life.

### $\tau(p \rightarrow e^+ \eta)$

**T4**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;4200</b>	<b><i>p</i></b>	<b>90</b>	<b>0</b>	<b>0.44</b>	NISHINO	12

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 81	<i>p</i>	90	1	1.7	WALL	00B	SOU2
> 313	<i>p</i>	90	0	0.2	MCGREW	99	IMB3
> 44	<i>p</i>	90	0	0.1	BERGER	91	FREJ
> 140	<i>p</i>	90	0	<0.04	HIRATA	89C	KAMI
> 100	<i>p</i>	90	0	0.6	SEIDEL	88	IMB
> 200	<i>p</i>	90	5	3.3	HAINES	86	IMB
> 64	<i>p</i>	90	0	<0.8	ARISAKA	85	KAMI
> 64	<i>p</i> (free)	90	5	6.5	BLEWITT	85	IMB
> 200	<i>p</i>	90	5	4.7	BLEWITT	85	IMB
> 1.2	<i>p</i>	90	2		<sup>1</sup> CHERRY	81	HOME

<sup>1</sup> We have converted 2 possible events to 90% CL limit.

### $\tau(p \rightarrow \mu^+ \eta)$

**T5**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;1300</b>	<b><i>p</i></b>	<b>90</b>	<b>2</b>	<b>0.49</b>	NISHINO	12

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 89	<i>p</i>	90	0	1.6	WALL	00B	SOU2
> 126	<i>p</i>	90	3	2.8	MCGREW	99	IMB3
> 26	<i>p</i>	90	1	0.8	BERGER	91	FREJ
> 69	<i>p</i>	90	1	<0.08	HIRATA	89C	KAMI
> 1.3	<i>p</i>	90	0	0.7	PHILLIPS	89	HPW
> 34	<i>p</i>	90	1	1.5	SEIDEL	88	IMB
> 46	<i>p</i>	90	7	6	HAINES	86	IMB
> 26	<i>p</i>	90	1	<0.8	ARISAKA	85	KAMI
> 17	<i>p</i> (free)	90	6	6	BLEWITT	85	IMB
> 46	<i>p</i>	90	7	8	BLEWITT	85	IMB

### $\tau(n \rightarrow \nu\eta)$

$\tau_6$

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN	
<b>&gt;158</b>	<i>n</i>	<b>90</b>	<b>0</b>	<b>1.2</b>	MCGREW	99	IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 71	<i>n</i>	90	2	3.7	WALL	00B	SOU2
> 29	<i>n</i>	90	0	0.9	BERGER	89	FREJ
> 54	<i>n</i>	90	2	0.9	HIRATA	89C	KAMI
> 16	<i>n</i>	90	3	2.1	SEIDEL	88	IMB
> 25	<i>n</i>	90	7	6	HAINES	86	IMB
> 30	<i>n</i>	90	0	0.4	KAJITA	86	KAMI
> 18	<i>n</i>	90	4	3	PARK	85	IMB
> 0.6	<i>n</i>	90	2		<sup>1</sup> CHERRY	81	HOME

<sup>1</sup> We have converted 2 possible events to 90% CL limit.

### $\tau(N \rightarrow e^+ \rho)$

$\tau_7$

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN	
<b>&gt;710</b>	<i>p</i>	<b>90</b>	<b>0</b>	<b>0.35</b>	NISHINO	12	SKAM
<b>&gt;217</b>	<i>n</i>	<b>90</b>	<b>4</b>	<b>4.8</b>	MCGREW	99	IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 70	<i>n</i>	90	1	0.38	NISHINO	12	SKAM
> 29	<i>p</i>	90	0	2.2	BERGER	91	FREJ
> 41	<i>n</i>	90	0	1.4	BERGER	91	FREJ
> 75	<i>p</i>	90	2	2.7	HIRATA	89C	KAMI
> 58	<i>n</i>	90	0	1.9	HIRATA	89C	KAMI
> 38	<i>n</i>	90	2	4.1	SEIDEL	88	IMB
> 1.2	<i>p</i>	90	0		BARTEL	87	SOUD
> 1.5	<i>n</i>	90	0		BARTEL	87	SOUD
> 17	<i>p</i>	90	7	7	HAINES	86	IMB
> 14	<i>n</i>	90	9	4	HAINES	86	IMB
> 12	<i>p</i>	90	0	<1.2	ARISAKA	85	KAMI
> 6	<i>n</i>	90	2	<1	ARISAKA	85	KAMI
> 6.7	<i>p</i> (free)	90	6	6	BLEWITT	85	IMB
> 17	<i>p</i>	90	7	7	BLEWITT	85	IMB
> 12	<i>n</i>	90	4	2	PARK	85	IMB

> 0.6	<i>n</i>	90	1	0.3	<sup>1</sup> BARTEL T	83	Soud
> 0.5	<i>p</i>	90	1	0.3	<sup>1</sup> BARTEL T	83	Soud
> 9.8	<i>p</i>	90	1		<sup>2</sup> KRISHNA...	82	KOLR
> 0.8	<i>p</i>	90	2		<sup>3</sup> CHERRY	81	HOME

<sup>1</sup> Limit based on zero events.<sup>2</sup> We have calculated 90% CL limit from 0 confined events.<sup>3</sup> We have converted 2 possible events to 90% CL limit. **$\tau(N \rightarrow \mu^+ \rho)$**  **$\tau_8$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN	
<b>&gt;160</b>	<i>p</i>	<b>90</b>	<b>1</b>	<b>0.42</b>	NISHINO	12	SKAM
<b>&gt;228</b>	<i>n</i>	<b>90</b>	<b>3</b>	<b>9.5</b>	MCGREW	99	IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 36	<i>n</i>	90	0	0.29	NISHINO	12	SKAM
> 12	<i>p</i>	90	0	0.5	BERGER	91	FREJ
> 22	<i>n</i>	90	0	1.1	BERGER	91	FREJ
>110	<i>p</i>	90	0	1.7	HIRATA	89C	KAMI
> 23	<i>n</i>	90	1	1.8	HIRATA	89C	KAMI
> 4.3	<i>p</i>	90	0	0.7	PHILLIPS	89	HPW
> 30	<i>p</i>	90	0	0.5	SEIDEL	88	IMB
> 11	<i>n</i>	90	1	1.1	SEIDEL	88	IMB
> 16	<i>p</i>	90	4	4.5	HAINES	86	IMB
> 7	<i>n</i>	90	6	5	HAINES	86	IMB
> 12	<i>p</i>	90	0	<0.7	ARISAKA	85	KAMI
> 5	<i>n</i>	90	1	<1.2	ARISAKA	85	KAMI
> 5.5	<i>p</i> (free)	90	4	5	BLEWITT	85	IMB
> 16	<i>p</i>	90	4	5	BLEWITT	85	IMB
> 9	<i>n</i>	90	1	2	PARK	85	IMB

 **$\tau(N \rightarrow \nu \rho)$**  **$\tau_9$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN	
<b>&gt;162</b>	<i>p</i>	<b>90</b>	<b>18</b>	<b>21.7</b>	MCGREW	99	IMB3
<b>&gt; 19</b>	<i>n</i>	<b>90</b>	<b>0</b>	<b>0.5</b>	SEIDEL	88	IMB

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 9	<i>n</i>	90	4	2.4	BERGER	89	FREJ
> 24	<i>p</i>	90	0	0.9	BERGER	89	FREJ
> 27	<i>p</i>	90	5	1.5	HIRATA	89C	KAMI
> 13	<i>n</i>	90	4	3.6	HIRATA	89C	KAMI
> 13	<i>p</i>	90	1	1.1	SEIDEL	88	IMB
> 8	<i>p</i>	90	6	5	HAINES	86	IMB
> 2	<i>n</i>	90	15	10	HAINES	86	IMB
> 11	<i>p</i>	90	2	1	KAJITA	86	KAMI
> 4	<i>n</i>	90	2	2	KAJITA	86	KAMI
> 4.1	<i>p</i> (free)	90	6	7	BLEWITT	85	IMB
> 8.4	<i>p</i>	90	6	5	BLEWITT	85	IMB
> 2	<i>n</i>	90	7	3	PARK	85	IMB
> 0.9	<i>p</i>	90	2		<sup>1</sup> CHERRY	81	HOME
> 0.6	<i>n</i>	90	2		<sup>1</sup> CHERRY	81	HOME

<sup>1</sup> We have converted 2 possible events to 90% CL limit.

$\tau(p \rightarrow e^+ \omega)$  $\tau_{10}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;320</b>	<b>p</b>	<b>90</b>	<b>1</b>	<b>0.53</b>	NISHINO	12
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>107	p	90	7	10.8	MCGREW	99
> 17	p	90	0	1.1	BERGER	91
> 45	p	90	2	1.45	HIRATA	89C
> 26	p	90	1	1.0	SEIDEL	88
> 1.5	p	90	0		BARTEL	87
> 37	p	90	6	5.3	HAINES	86
> 25	p	90	1	<1.4	ARISAKA	85
> 12	p (free)	90	6	7.5	BLEWITT	85
> 37	p	90	6	5.7	BLEWITT	85
> 0.6	p	90	1	0.3	<sup>1</sup> BARTEL	83
> 9.8	p	90	1		<sup>2</sup> KRISHNA...	82
> 2.8	p	90	2		<sup>3</sup> CHERRY	81
					HOME	

<sup>1</sup> Limit based on zero events.<sup>2</sup> We have calculated 90% CL limit from 0 confined events.<sup>3</sup> We have converted 2 possible events to 90% CL limit. $\tau(p \rightarrow \mu^+ \omega)$  $\tau_{11}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;780</b>	<b>p</b>	<b>90</b>	<b>0</b>	<b>0.48</b>	NISHINO	12
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>117	p	90	11	12.1	MCGREW	99
> 11	p	90	0	1.0	BERGER	91
> 57	p	90	2	1.9	HIRATA	89C
> 4.4	p	90	0	0.7	PHILLIPS	89
> 10	p	90	2	1.3	SEIDEL	88
> 23	p	90	2	1	HAINES	86
> 6.5	p (free)	90	9	8.7	BLEWITT	85
> 23	p	90	8	7	BLEWITT	85
					IMB	

 $\tau(n \rightarrow \nu \omega)$  $\tau_{12}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;108</b>	<b>n</b>	<b>90</b>	<b>12</b>	<b>22.5</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
> 17	n	90	1	0.7	BERGER	89
> 43	n	90	3	2.7	HIRATA	89C
> 6	n	90	2	1.3	SEIDEL	88
> 12	n	90	6	6	HAINES	86
> 18	n	90	2	2	KAJITA	86
> 16	n	90	1	2	PARK	85
> 2.0	n	90	2		<sup>1</sup> CHERRY	81
					HOME	

<sup>1</sup> We have converted 2 possible events to 90% CL limit.

$\tau(N \rightarrow e^+ K)$  $\tau_{13}$ 

<u>LIMIT</u> ( $10^{-30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
> 1000	p	90	6	4.7	KOBAYASHI	05
> 17	n	90	35	29.4	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 85	p	90	3	4.9	WALL	00	SOU2
> 31	p	90	23	25.2	MCGREW	99	IMB3
> 60	p	90	0		BERGER	91	FREJ
> 150	p	90	0	<0.27	HIRATA	89C	KAMI
> 70	p	90	0	1.8	SEIDEL	88	IMB
> 77	p	90	5	4.5	HAINES	86	IMB
> 38	p	90	0	<0.8	ARISAKA	85	KAMI
> 24	p (free)	90	7	8.5	BLEWITT	85	IMB
> 77	p	90	5	4	BLEWITT	85	IMB
> 1.3	p	90	0		ALEKSEEV	81	BAKS
> 1.3	n	90	0		ALEKSEEV	81	BAKS

 $\tau(p \rightarrow e^+ K_S^0)$  $\tau_{14}$ 

<u>LIMIT</u> ( $10^{-30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
> 120	p	90	1	1.3	WALL	00
> 76	p	90	0	0.5	BERGER	91

 $\tau(p \rightarrow e^+ K_L^0)$  $\tau_{15}$ 

<u>LIMIT</u> ( $10^{-30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
> 51	p	90	2	3.5	WALL	00
> 44	p	90	0	$\leq 0.1$	BERGER	91

 $\tau(N \rightarrow \mu^+ K)$  $\tau_{16}$ 

<u>LIMIT</u> ( $10^{-30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
> 1600	p	90	13	13.2	REGIS	12
> 26	n	90	20	28.4	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 1300	p	90	3	3.9	KOBAYASHI	05	SKAM
> 120	p	90	0	<1.2	WALL	00	SOU2
> 120	p	90	4	7.2	MCGREW	99	IMB3
> 54	p	90	0		BERGER	91	FREJ
> 120	p	90	1	0.4	HIRATA	89C	KAMI
> 3.0	p	90	0	0.7	PHILLIPS	89	HPW
> 19	p	90	3	2.5	SEIDEL	88	IMB
> 1.5	p	90	0		BARTEL	87	SOUD
> 1.1	n	90	0		BARTEL	87	SOUD
> 40	p	90	7	6	HAINES	86	IMB

> 19	$p$	90	1 <1.1	ARISAKA	85	KAMI
> 6.7	$p$ (free)	90	11 13	BLEWITT	85	IMB
> 40	$p$	90	7 8	BLEWITT	85	IMB
> 6	$p$	90	1	BATTISTONI	84	NUSX
> 0.6	$p$	90	0	<sup>2</sup> BARTEL	83	SOUD
> 0.4	$n$	90	0	<sup>2</sup> BARTEL	83	SOUD
> 5.8	$p$	90	2	<sup>3</sup> KRISHNA...	82	KOLR
> 2.0	$p$	90	0	CHERRY	81	HOME
> 0.2	$n$	90		<sup>4</sup> GURR	67	CNTR

<sup>1</sup> BARTEL 87 limit applies to  $p \rightarrow \mu^+ K_S^0$ .

<sup>2</sup> Limit based on zero events.

<sup>3</sup> We have calculated 90% CL limit from 1 confined event.

<sup>4</sup> We have converted half-life to 90% CL mean life.

### $\tau(p \rightarrow \mu^+ K_S^0)$

**$\tau_{17}$**

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>150	$p$	90	0 <0.8		WALL	00 SOU2
> 64	$p$	90	0 1.2		BERGER	91 FREJ

### $\tau(p \rightarrow \mu^+ K_L^0)$

**$\tau_{18}$**

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>83	$p$	90	0 0.4		WALL	00 SOU2
>44	$p$	90	0 $\leq 0.1$		BERGER	91 FREJ

### $\tau(N \rightarrow \nu K)$

**$\tau_{19}$**

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;5900</b>	<b><math>p</math></b>	<b>90</b>	<b>0</b>	<b>1.0</b>	ABE	14G SKAM
<b>&gt; 86</b>	<b><math>n</math></b>	<b>90</b>	<b>0</b>	<b>2.4</b>	HIRATA	89C KAMI
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 540	$p$	90	0 0.9		ASAKURA	15 KLND
>2300	$p$	90	0 1.3		KOBAYASHI	05 SKAM
> 26	$n$	90	16 9.1		WALL	00 SOU2
> 670	$p$	90			HAYATO	99 SKAM
> 151	$p$	90	15 21.4		MCGREW	99 IMB3
> 30	$n$	90	34 34.1		MCGREW	99 IMB3
> 43	$p$	90	1 1.54		<sup>1</sup> ALLISON	98 SOU2
> 15	$n$	90	1 1.8		BERGER	89 FREJ
> 15	$p$	90	1 1.8		BERGER	89 FREJ
> 100	$p$	90	9 7.3		HIRATA	89C KAMI
> 0.28	$p$	90	0 0.7		PHILLIPS	89 HPW
> 0.3	$p$	90	0		BARTEL	87 SOUD
> 0.75	$n$	90	0		<sup>2</sup> BARTEL	87 SOUD
> 10	$p$	90	6 5		HAINES	86 IMB

> 15	<i>n</i>	90	3 5	HAINES	86	IMB
> 28	<i>p</i>	90	3 3	KAJITA	86	KAMI
> 32	<i>n</i>	90	0 1.4	KAJITA	86	KAMI
> 1.8	<i>p</i> (free)	90	6 11	BLEWITT	85	IMB
> 9.6	<i>p</i>	90	6 5	BLEWITT	85	IMB
> 10	<i>n</i>	90	2 2	PARK	85	IMB
> 5	<i>n</i>	90	0	BATTISTONI	84	NUSX
> 2	<i>p</i>	90	0	BATTISTONI	84	NUSX
> 0.3	<i>n</i>	90	0	<sup>3</sup> BARTEL	83	SOU
> 0.1	<i>p</i>	90	0	<sup>3</sup> BARTEL	83	SOU
> 5.8	<i>p</i>	90	1	<sup>4</sup> KRISHNA...	82	KOLR
> 0.3	<i>n</i>	90	2	<sup>5</sup> CHERRY	81	HOME

<sup>1</sup> This ALLISON 98 limit is with no background subtraction; with subtraction the limit becomes  $> 46 \times 10^{30}$  years.

<sup>2</sup> BARTEL 87 limit applies to  $n \rightarrow \nu K_S^0$ .

<sup>3</sup> Limit based on zero events.

<sup>4</sup> We have calculated 90% CL limit from 1 confined event.

<sup>5</sup> We have converted 2 possible events to 90% CL limit.

### $\tau(n \rightarrow \nu K_S^0)$

**$\tau_{20}$**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>260	<i>n</i>	90	34	30	<sup>1</sup> KOBAYASHI 05	SKAM

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 51	<i>n</i>	90	16	9.1	WALL	00	SOU2
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<sup>1</sup> We have doubled the  $n \rightarrow \nu K^0$  limit given in KOBAYASHI 05 to obtain this  $n \rightarrow \nu K_S^0$  limit.

### $\tau(p \rightarrow e^+ K^*(892)^0)$

**$\tau_{21}$**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>84	<i>p</i>	90	38	52.0	MCGREW 99	IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

>10	<i>p</i>	90	0	0.8	BERGER	91	FREJ
>52	<i>p</i>	90	2	1.55	HIRATA	89C	KAMI
>10	<i>p</i>	90	1	<1	ARISAKA	85	KAMI

### $\tau(N \rightarrow \nu K^*(892)^0)$

**$\tau_{22}$**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>51	<i>p</i>	90	7	9.1	MCGREW 99	IMB3
>78	<i>n</i>	90	40	50	MCGREW 99	IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

>22	<i>n</i>	90	0	2.1	BERGER	89	FREJ
>17	<i>p</i>	90	0	2.4	BERGER	89	FREJ
>20	<i>p</i>	90	5	2.1	HIRATA	89C	KAMI
>21	<i>n</i>	90	4	2.4	HIRATA	89C	KAMI
>10	<i>p</i>	90	7	6	HAINES	86	IMB
> 5	<i>n</i>	90	8	7	HAINES	86	IMB

> 8	<i>p</i>	90	3 2	KAJITA	86	KAMI
> 6	<i>n</i>	90	2 1.6	KAJITA	86	KAMI
> 5.8	<i>p</i> (free)	90	10 16	BLEWITT	85	IMB
> 9.6	<i>p</i>	90	7 6	BLEWITT	85	IMB
> 7	<i>n</i>	90	1 4	PARK	85	IMB
> 2.1	<i>p</i>	90	1	<sup>1</sup> BATTISTONI	82	NUSX

<sup>1</sup> We have converted 1 possible event to 90% CL limit.

### — Antilepton + mesons —

#### $\tau(p \rightarrow e^+ \pi^+ \pi^-)$

**T23**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;82</b>	<b><i>p</i></b>	<b>90</b>	<b>16</b>	<b>23.1</b>	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

>21	<i>p</i>	90	0	2.2	BERGER	91	FREJ
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#### $\tau(p \rightarrow e^+ \pi^0 \pi^0)$

**T24**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;147</b>	<b><i>p</i></b>	<b>90</b>	<b>2</b>	<b>0.8</b>	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 38	<i>p</i>	90	1	0.5	BERGER	91	FREJ
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#### $\tau(n \rightarrow e^+ \pi^- \pi^0)$

**T25**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;52</b>	<b><i>n</i></b>	<b>90</b>	<b>38</b>	<b>34.2</b>	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

>32	<i>n</i>	90	1	0.8	BERGER	91	FREJ
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#### $\tau(p \rightarrow \mu^+ \pi^+ \pi^-)$

**T26**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;133</b>	<b><i>p</i></b>	<b>90</b>	<b>25</b>	<b>38.0</b>	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 17	<i>p</i>	90	1	2.6	BERGER	91	FREJ
> 3.3	<i>p</i>	90	0	0.7	PHILLIPS	89	HPW

#### $\tau(p \rightarrow \mu^+ \pi^0 \pi^0)$

**T27**

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
<b>&gt;101</b>	<b><i>p</i></b>	<b>90</b>	<b>3</b>	<b>1.6</b>	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 33	<i>p</i>	90	1	0.9	BERGER	91	FREJ
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$\tau(n \rightarrow \mu^+ \pi^- \pi^0)$  $\tau_{28}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;74</b>	<b><i>n</i></b>	<b>90</b>	<b>17</b>	<b>20.8</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>33	<i>n</i>	90	0	1.1	BERGER	91
					FREJ	

 $\tau(n \rightarrow e^+ K^0 \pi^-)$  $\tau_{29}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;18</b>	<b><i>n</i></b>	<b>90</b>	<b>1</b>	<b>0.2</b>	BERGER	91
					FREJ	

**Lepton + meson** $\tau(n \rightarrow e^- \pi^+)$  $\tau_{30}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;65</b>	<b><i>n</i></b>	<b>90</b>	<b>0</b>	<b>1.6</b>	SEIDEL	88
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>55	<i>n</i>	90	0	1.09	BERGER	91B
>16	<i>n</i>	90	9	7	HAINES	86
>25	<i>n</i>	90	2	4	PARK	85
					IMB	

 $\tau(n \rightarrow \mu^- \pi^+)$  $\tau_{31}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;49</b>	<b><i>n</i></b>	<b>90</b>	<b>0</b>	<b>0.5</b>	SEIDEL	88
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>33	<i>n</i>	90	0	1.40	BERGER	91B
> 2.7	<i>n</i>	90	0	0.7	PHILLIPS	89
>25	<i>n</i>	90	7	6	HAINES	86
>27	<i>n</i>	90	2	3	PARK	85
					IMB	

 $\tau(n \rightarrow e^- \rho^+)$  $\tau_{32}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;62</b>	<b><i>n</i></b>	<b>90</b>	<b>2</b>	<b>4.1</b>	SEIDEL	88
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>12	<i>n</i>	90	13	6	HAINES	86
>12	<i>n</i>	90	5	3	PARK	85
					IMB	

 $\tau(n \rightarrow \mu^- \rho^+)$  $\tau_{33}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;7</b>	<b><i>n</i></b>	<b>90</b>	<b>1</b>	<b>1.1</b>	SEIDEL	88
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>2.6	<i>n</i>	90	0	0.7	PHILLIPS	89
>9	<i>n</i>	90	7	5	HAINES	86
>9	<i>n</i>	90	2	2	PARK	85
					IMB	

$\tau(n \rightarrow e^- K^+)$  $\tau_{34}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;32</b>	<b>n</b>	<b>90</b>	<b>3</b>	<b>2.96</b>	BERGER	91B FREJ

• • • We do not use the following data for averages, fits, limits, etc. • • •

&gt; 0.23      n      90      0    0.7

PHILLIPS      89    HPW

 $\tau(n \rightarrow \mu^- K^+)$  $\tau_{35}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;57</b>	<b>n</b>	<b>90</b>	<b>0</b>	<b>2.18</b>	BERGER	91B FREJ

• • • We do not use the following data for averages, fits, limits, etc. • • •

&gt; 4.7      n      90      0    0.7

PHILLIPS      89    HPW

**Lepton + mesons** $\tau(p \rightarrow e^- \pi^+ \pi^+)$  $\tau_{36}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;30</b>	<b>p</b>	<b>90</b>	<b>1</b>	<b>2.50</b>	BERGER	91B FREJ

• • • We do not use the following data for averages, fits, limits, etc. • • •

&gt; 2.0      p      90      0    0.7

PHILLIPS      89    HPW

 $\tau(n \rightarrow e^- \pi^+ \pi^0)$  $\tau_{37}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;29</b>	<b>n</b>	<b>90</b>	<b>1</b>	<b>0.78</b>	BERGER	91B FREJ

 $\tau(p \rightarrow \mu^- \pi^+ \pi^+)$  $\tau_{38}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;17</b>	<b>p</b>	<b>90</b>	<b>1</b>	<b>1.72</b>	BERGER	91B FREJ

• • • We do not use the following data for averages, fits, limits, etc. • • •

&gt; 7.8      p      90      0    0.7

PHILLIPS      89    HPW

 $\tau(n \rightarrow \mu^- \pi^+ \pi^0)$  $\tau_{39}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;34</b>	<b>n</b>	<b>90</b>	<b>0</b>	<b>0.78</b>	BERGER	91B FREJ

 $\tau(p \rightarrow e^- \pi^+ K^+)$  $\tau_{40}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;75</b>	<b>p</b>	<b>90</b>	<b>81</b>	<b>127.2</b>	MCGREW	99 IMB3

• • • We do not use the following data for averages, fits, limits, etc. • • •

&gt;20      p      90      3    2.50

BERGER      91B FREJ

$\tau(p \rightarrow \mu^- \pi^+ K^+)$ **T41**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>245	<i>p</i>	90	3	4.0	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
> 5	<i>p</i>	90	2	0.78	BERGER	91B
<b>— Antilepton + photon(s) —</b>						

 $\tau(p \rightarrow e^+ \gamma)$ **T42**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>670	<i>p</i>	90	0	0.1	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>133	<i>p</i>	90	0	0.3	BERGER	91
>460	<i>p</i>	90	0	0.6	SEIDEL	88
>360	<i>p</i>	90	0	0.3	HAINES	86
> 87	<i>p</i> (free)	90	0	0.2	BLEWITT	85
>360	<i>p</i>	90	0	0.2	BLEWITT	85
> 0.1	<i>p</i>	90			<sup>1</sup> GURR	67
<sup>1</sup> We have converted half-life to 90% CL mean life.						

 $\tau(p \rightarrow \mu^+ \gamma)$ **T43**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>478	<i>p</i>	90	0	0.1	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>155	<i>p</i>	90	0	0.1	BERGER	91
>380	<i>p</i>	90	0	0.5	SEIDEL	88
> 97	<i>p</i>	90	3	2	HAINES	86
> 61	<i>p</i> (free)	90	0	0.2	BLEWITT	85
>280	<i>p</i>	90	0	0.6	BLEWITT	85
> 0.3	<i>p</i>	90			<sup>1</sup> GURR	67
<sup>1</sup> We have converted half-life to 90% CL mean life.						

 $\tau(n \rightarrow \nu \gamma)$ **T44**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>550	<i>n</i>	90			TAKHISTOV	15
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
> 28	<i>n</i>	90	163	144.7	MCGREW	99
> 24	<i>n</i>	90	10	6.86	BERGER	91B
> 9	<i>n</i>	90	73	60	HAINES	86
> 11	<i>n</i>	90	28	19	PARK	85

 $\tau(p \rightarrow e^+ \gamma\gamma)$ **T45**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>100	<i>p</i>	90	1	0.8	BERGER	91

$\tau(n \rightarrow \nu\gamma\gamma)$ **T46**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>219	<i>n</i>	90	5	7.5	MCGREW	99

**Antilepton + single massless** $\tau(p \rightarrow e^+ X)$ **T47**

<i>VALUE</i> ( $10^{30}$ years)	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>790	90	TAKHISTOV	15

 $\tau(p \rightarrow \mu^+ X)$ **T48**

<i>VALUE</i> ( $10^{30}$ years)	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>410	90	TAKHISTOV	15

**Three (or more) leptons** $\tau(p \rightarrow e^+ e^+ e^-)$ **T49**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>793	<i>p</i>	90	0	0.5	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

>147	<i>p</i>	90	0	0.1	BERGER	91	FREJ
>510	<i>p</i>	90	0	0.3	HAINES	86	IMB
> 89	<i>p</i> (free)	90	0	0.5	BLEWITT	85	IMB
>510	<i>p</i>	90	0	0.7	BLEWITT	85	IMB

 $\tau(p \rightarrow e^+ \mu^+ \mu^-)$ **T50**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>359	<i>p</i>	90	1	0.9	MCGREW	99

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 81	<i>p</i>	90	0	0.16	BERGER	91	FREJ
> 5.0	<i>p</i>	90	0	0.7	PHILLIPS	89	HPW

 $\tau(p \rightarrow e^+ \nu\nu)$ **T51**

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
>170	<i>p</i>	90			<sup>1</sup> TAKHISTOV	14

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 17	<i>p</i>	90	152	153.7	MCGREW	99	IMB3
> 11	<i>p</i>	90	11	6.08	BERGER	91B	FREJ

1 Allowed events at 90% CL are 459.

$\tau(n \rightarrow e^+ e^- \nu)$  $\tau_{52}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;257</b>	<b>n</b>	<b>90</b>	<b>5</b>	<b>7.5</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
> 74	n	90	0	< 0.1	BERGER	91B
> 45	n	90	5	5	HAINES	86
> 26	n	90	4	3	PARK	85

 $\tau(n \rightarrow \mu^+ \mu^- \nu)$  $\tau_{53}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;83</b>	<b>n</b>	<b>90</b>	<b>25</b>	<b>29.4</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>47	n	90	0	< 0.1	BERGER	91B

 $\tau(n \rightarrow \mu^+ \mu^- \nu)$  $\tau_{54}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;79</b>	<b>n</b>	<b>90</b>	<b>100</b>	<b>145</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>42	n	90	0	1.4	BERGER	91B
> 5.1	n	90	0	0.7	PHILLIPS	89
>16	n	90	14	7	HAINES	86
>19	n	90	4	7	PARK	85

 $\tau(p \rightarrow \mu^+ e^+ e^-)$  $\tau_{55}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;529</b>	<b>p</b>	<b>90</b>	<b>0</b>	<b>1.0</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
> 91	p	90	0	$\leq 0.1$	BERGER	91

 $\tau(p \rightarrow \mu^+ \mu^+ \mu^-)$  $\tau_{56}$ 

<u>LIMIT</u> ( $10^{30}$ years)	<u>PARTICLE</u>	<u>CL%</u>	<u>EVTS</u>	<u>BKGD EST</u>	<u>DOCUMENT ID</u>	<u>TECN</u>
<b>&gt;675</b>	<b>p</b>	<b>90</b>	<b>0</b>	<b>0.3</b>	MCGREW	99
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>						
>119	p	90	0	0.2	BERGER	91
> 10.5	p	90	0	0.7	PHILLIPS	89
>190	p	90	1	0.1	HAINES	86
> 44	p (free)	90	1	0.7	BLEWITT	85
>190	p	90	1	0.9	BLEWITT	85
> 2.1	p	90	1		<sup>1</sup> BATTISTONI	82

<sup>1</sup> We have converted 1 possible event to 90% CL limit.

$\tau(p \rightarrow \mu^+ \nu \nu)$  **$\tau_{57}$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>220	p	90			1 TAKHISTOV	14 SKAM

• • • We do not use the following data for averages, fits, limits, etc. • • •

> 21	p	90	7	11.23	BERGER	91B FREJ
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<sup>1</sup> Allowed events at 90% CL are 286.

 $\tau(p \rightarrow e^- \mu^+ \mu^+)$  **$\tau_{58}$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>6.0	p	90	0	0.7	PHILLIPS	89 HPW

 $\tau(n \rightarrow 3\nu)$  **$\tau_{59}$** 

See also the “to anything” and “disappearance” limits for bound nucleons in the “*p* Mean Life” data block just in front of the list of possible *p* decay modes. Such modes could of course be to three (or five) neutrinos, and the limits are stronger, but we do not repeat them here.

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>0.00049	n	90	2	2	1 SUZUKI	93B KAMI

• • • We do not use the following data for averages, fits, limits, etc. • • •

>0.0023	n	90			2 GLICENSTEIN	97 KAMI
>0.00003	n	90	11	6.1	3 BERGER	91B FREJ
>0.00012	n	90	7	11.2	3 BERGER	91B FREJ
>0.0005	n	90	0		LEARNED	79 RVUE

<sup>1</sup> The SUZUKI 93B limit applies to any of  $\nu_e \nu_e \bar{\nu}_e$ ,  $\nu_\mu \nu_\mu \bar{\nu}_\mu$ , or  $\nu_\tau \nu_\tau \bar{\nu}_\tau$ .

<sup>2</sup> GLICENSTEIN 97 uses Kamioka data and the idea that the disappearance of the neutron's magnetic moment should produce radiation.

<sup>3</sup> The first BERGER 91B limit is for  $n \rightarrow \nu_e \nu_e \bar{\nu}_e$ , the second is for  $n \rightarrow \nu_\mu \nu_\mu \bar{\nu}_\mu$ .

 $\tau(n \rightarrow 5\nu)$  **$\tau_{60}$** 

See the note on  $\tau(n \rightarrow 3\nu)$  on the previous data block.

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>0.0017	n	90			1 GLICENSTEIN	97 KAMI

<sup>1</sup> GLICENSTEIN 97 uses Kamioka data and the idea that the disappearance of the neutron's magnetic moment should produce radiation.

**Inclusive modes** $\tau(N \rightarrow e^+ \text{anything})$  **$\tau_{61}$** 

LIMIT ( $10^{30}$ years)	PARTICLE	CL%	EVTS	BKGD EST	DOCUMENT ID	TECN
>0.6	p, n	90			1 LEARNED	79 RVUE

<sup>1</sup> The electron may be primary or secondary.

$\tau(N \rightarrow \mu^+ \text{anything})$  $\tau_{62}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;12</b>	<b>p, n</b>	<b>90</b>	<b>2</b>		1,2 CHERRY	81 HOME
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 1.8	<i>p, n</i>	90			<sup>2</sup> COWSIK	80 CNTR
> 6	<i>p, n</i>	90			<sup>2</sup> LEARNED	79 RVUE

<sup>1</sup> We have converted 2 possible events to 90% CL limit.<sup>2</sup> The muon may be primary or secondary. $\tau(N \rightarrow \nu \text{anything})$  $\tau_{63}$ Anything =  $\pi$ ,  $\rho$ ,  $K$ , etc.

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>0.0002	<i>p, n</i>	90	0		LEARNED	79 RVUE

 $\tau(N \rightarrow e^+ \pi^0 \text{anything})$  $\tau_{64}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;0.6</b>	<b>p, n</b>	<b>90</b>	<b>0</b>		LEARNED	79 RVUE

 $\tau(N \rightarrow 2 \text{ bodies}, \nu\text{-free})$  $\tau_{65}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>PARTICLE</i>	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>1.3	<i>p, n</i>	90	0		ALEKSEEV	81 BAKS

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 $\Delta B = 2$  dinucleon modes

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 $\tau(pp \rightarrow \pi^+ \pi^+)$  $\tau_{66}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;72.2</b>	<b>90</b>	<b>2</b>	<b>4.45</b>	GUSTAFSON	15 SKAM	per oxygen nucleus
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 0.7	90	4	2.34	BERGER	91B FREJ	per iron nucleus

 $\tau(pn \rightarrow \pi^+ \pi^0)$  $\tau_{67}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;170</b>	<b>90</b>			GUSTAFSON	15 SKAM	per oxygen nucleus
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 2.0	90	0	0.31	BERGER	91B FREJ	per iron nucleus

 $\tau(nn \rightarrow \pi^+ \pi^-)$  $\tau_{68}$ 

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;0.7</b>	<b>90</b>	<b>4</b>	<b>2.18</b>	BERGER	91B FREJ	$\tau$ per iron nucleus

$\tau(nn \rightarrow \pi^0\pi^0)$ **T69**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;404</b>	<b>90</b>			GUSTAFSON	15	SKAM per oxygen nucleus
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 3.4	90	0	0.78	BERGER	91B	FREJ per iron nucleus

 $\tau(pp \rightarrow K^+K^+)$ **T70**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;170</b>	<b>90</b>	<b>0</b>	<b>0.28</b>	LITOS	14	SKAM $\tau$ per oxygen nucleus

 $\tau(pp \rightarrow e^+e^+)$ **T71**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;5.8</b>	<b>90</b>	<b>0</b>	<b>&lt;0.1</b>	BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pp \rightarrow e^+\mu^+)$ **T72**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;3.6</b>	<b>90</b>	<b>0</b>	<b>&lt;0.1</b>	BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pp \rightarrow \mu^+\mu^+)$ **T73**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;1.7</b>	<b>90</b>	<b>0</b>	<b>0.62</b>	BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pn \rightarrow e^+\bar{\nu})$ **T74**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;260</b>	<b>90</b>			TAKHISTOV	15	SKAM
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 2.8	90	5	9.67	BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pn \rightarrow \mu^+\bar{\nu})$ **T75**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;200</b>	<b>90</b>			TAKHISTOV	15	SKAM
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 1.6	90	4	4.37	BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pn \rightarrow \tau^+\bar{\nu}_\tau)$ **T76**

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;29</b>	<b>90</b>			TAKHISTOV	15	SKAM
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
> 1	90			<sup>1</sup> BRYMAN	14	CHER

<sup>1</sup> BRYMAN 14 uses a MCGREW 99 limit on the  $p \rightarrow e^+\nu\nu$  lifetime to extract this value.

$\tau(nn \rightarrow \nu_e \bar{\nu}_e)$ **T77**

We include “invisible” modes here.

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;1.4</b>	<b>90</b>			<sup>1</sup> ARAKI	06	KLND $nn \rightarrow$ invisible
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>0.000042 90				<sup>2</sup> TRETYAK	04	CNTR $nn \rightarrow$ invisible
>0.000049 90				<sup>3</sup> BACK	03	BORX $nn \rightarrow$ invisible
>0.000012 90				<sup>4</sup> BERNABEI	00B	DAMA $nn \rightarrow$ invisible
>0.000012 90	5	9.7		BERGER	91B	FREJ $\tau$ per iron nucleus

<sup>1</sup> ARAKI 06 looks for signs of de-excitation of the residual nucleus after disappearance of two neutrons from the s shell of  $^{12}\text{C}$ .

<sup>2</sup> TRETYAK 04 uses data from an old Homestake-mine radiochemical experiment on limits for invisible decays of  $^{39}\text{K}$  to  $^{37}\text{Ar}$ .

<sup>3</sup> BACK 03 looks for decays of unstable nuclides left after  $NN$  decays of parent  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$  nuclei. These are “invisible channel” limits.

<sup>4</sup> BERNABEI 00B looks for the decay of a  $^{127}_{54}\text{Xe}$  nucleus following the disappearance of an  $nn$  pair in the otherwise-stable  $^{129}_{54}\text{Xe}$  nucleus. The limit here applies as well to  $nn \rightarrow \nu_\mu \bar{\nu}_\mu$ ,  $nn \rightarrow \nu_\tau \bar{\nu}_\tau$ , or any “disappearance” mode.

 $\tau(nn \rightarrow \nu_\mu \bar{\nu}_\mu)$ **T78**

See the proceeding data block. “Invisible modes” would include any multi-neutrino mode.

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
<b>&gt;1.4</b>	<b>(CL = 90%) OUR LIMIT</b>						
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>							
>0.000006 90	4	4.4			BERGER	91B	FREJ $\tau$ per iron nucleus

 $\tau(pn \rightarrow \text{invisible})$ **T79**

This violates charge conservation as well as baryon number conservation.

<i>VALUE</i> ( $10^{30}$ years)	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;0.000021</b>	90	<sup>1</sup> TRETYAK	04 CNTR

<sup>1</sup> TRETYAK 04 uses data from an old Homestake-mine radiochemical experiment on limits for invisible decays of  $^{39}\text{K}$  to  $^{37}\text{Ar}$ .

 $\tau(pp \rightarrow \text{invisible})$ **T80**

This violates charge conservation as well as baryon number conservation.

<i>LIMIT</i> ( $10^{30}$ years)	<i>CL%</i>	<i>EVTS</i>	<i>BKGD EST</i>	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>
<b>&gt;0.00005</b>		90		<sup>1</sup> BACK	03	BORX
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>						
>0.00000055 90				<sup>2</sup> BERNABEI	00B	DAMA

<sup>1</sup> BACK 03 looks for decays of unstable nuclides left after  $NN$  decays of parent  $^{12}\text{C}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$  nuclei. These are “invisible channel” limits.

<sup>2</sup> BERNABEI 00B looks for the decay of a  $^{127}_{52}\text{Te}$  nucleus following the disappearance of a  $pp$  pair in the otherwise-stable  $^{129}_{54}\text{Xe}$  nucleus.

## $\bar{p}$ PARTIAL MEAN LIVES

The “partial mean life” limits tabulated here are the limits on  $\bar{\tau}/B_i$ , where  $\bar{\tau}$  is the total mean life for the antiproton and  $B_i$  is the branching fraction for the mode in question.

### $\tau(\bar{p} \rightarrow e^- \gamma)$

**T81**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 7 \times 10^5$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
>1848	95	GEER 94	CALO	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow \mu^- \gamma)$

**T82**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 5 \times 10^4$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
$> 5.0 \times 10^4$	90	HU 98B	APEX	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow e^- \pi^0)$

**T83**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 4 \times 10^5$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
>554	95	GEER 94	CALO	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow \mu^- \pi^0)$

**T84**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 5 \times 10^4$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
$> 4.8 \times 10^4$	90	HU 98B	APEX	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow e^- \eta)$

**T85**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 2 \times 10^4$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
>171	95	GEER 94	CALO	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow \mu^- \eta)$

**T86**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 8 \times 10^3$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
$> 7.9 \times 10^3$	90	HU 98B	APEX	8.9 GeV/c $\bar{p}$ beam

### $\tau(\bar{p} \rightarrow e^- K_S^0)$

**T87**

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$> 900$	90	GEER 00	APEX	8.9 GeV/c $\bar{p}$ beam
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
> 29	95	GEER 94	CALO	8.9 GeV/c $\bar{p}$ beam

$\tau(\bar{p} \rightarrow \mu^- K_S^0)$  $\tau_{88}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>4 \times 10^3$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
$>4.3 \times 10^3$	90	HU	98B	APEX 8.9 GeV/c $\bar{p}$ beam

 $\tau(\bar{p} \rightarrow e^- K_L^0)$  $\tau_{89}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>9 \times 10^3$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
$>9$	95	GEER	94	CALO 8.9 GeV/c $\bar{p}$ beam

 $\tau(\bar{p} \rightarrow \mu^- K_L^0)$  $\tau_{90}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>7 \times 10^3$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
$>6.5 \times 10^3$	90	HU	98B	APEX 8.9 GeV/c $\bar{p}$ beam

 $\tau(\bar{p} \rightarrow e^- \gamma\gamma)$  $\tau_{91}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>2 \times 10^4$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam

 $\tau(\bar{p} \rightarrow \mu^- \gamma\gamma)$  $\tau_{92}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>2 \times 10^4$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
$>2.3 \times 10^4$	90	HU	98B	APEX 8.9 GeV/c $\bar{p}$ beam

 $\tau(\bar{p} \rightarrow e^- \omega)$  $\tau_{93}$ 

VALUE (years)	CL%	DOCUMENT ID	TECN	COMMENT
$>200$	90	GEER	00	APEX 8.9 GeV/c $\bar{p}$ beam

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LEARNED	79	PRL 43 907	J.G. Learned, F. Reines, A. Soni	(UCI)
BREGMAN	78	PL 78B 174	M. Bregman <i>et al.</i>	(CERN)
ROBERTS	78	PR D17 358	B.L. Roberts	(WILL, RHEL)
EVANS	77	SCI 197 989	J.C. Evans Jr., R.I. Steinberg	(BNL, PENN)
HU	75	NP A254 403	E. Hu <i>et al.</i>	(COLU, YALE)
COHEN	73	JPCRD 2 664	E.R. Cohen, B.N. Taylor	(RISC, NBS)
DYLLA	73	PR A7 1224	H.F. Dylla, J.G. King	(MIT)
DIX	70	Thesis Case	F.W. Dix	(CASE)
HARRISON	69	PRL 22 1263	G.E. Harrison, P.G.H. Sandars, S.J. Wright	(OXF)
GURR	67	PR 158 1321	H.S. Gurr <i>et al.</i>	(CASE, WITW)
FLOEROV	58	DOKL 3 79	G.N. Flerov <i>et al.</i>	(ASCI)